

```
In[111]:=
SetDirectory[NotebookDirectory[]];
<< MaTeX`
texStyle = {};
SetOptions[MaTeX,
  "BasePreamble" → {"\\usepackage{amsmath}", "\\usepackage{xcolor}",
    "\\usepackage{fourier}", "\\usepackage{ebgaramond}"}, FontSize → 11];
```

```
In[115]:=
frame[legend_] :=
  Framed[legend, FrameStyle → Thin, RoundingRadius → 10, FrameMargins → 0];
```

---

```
In[116]:=
Clear[findAllRoots]
(* from https://
  mathematica.stackexchange.com/questions/16439/find-all-roots-of-an-
  interpolating-function-solution-to-a-differential-equation/16444#16444 *)
SyntaxInformation[findAllRoots] = {"LocalVariables" → {"Plot", {2, 2}},
  "ArgumentsPattern" → {_, _, OptionsPattern[]}};
SetAttributes[findAllRoots, HoldAll];

Options[findAllRoots] = Join[{"ShowPlot" → False, PlotRange → All},
  FilterRules[Options[Plot], Except[PlotRange]]];

findAllRoots[fn_, {l_, lmin_, lmax_}, opts : OptionsPattern[]] :=
  Module[{pl, p, x, localFunction, brackets},
    localFunction = ReleaseHold[Hold[fn] /. HoldPattern[l] => x];
    If[lmin ≠ lmax, pl = Plot[localFunction, {x, lmin, lmax}, Evaluate@
      FilterRules[Join[{opts}, Options[findAllRoots]], Options[Plot]]];
    p = Cases[pl, Line[{x_}] => x, Infinity];
    If[OptionValue["ShowPlot"],
      Print[Show[pl, PlotLabel → "Finding roots for this function",
        ImageSize → 200, BaseStyle → {FontSize → 8}]]], p = {}];
    brackets =
      Map[First, Select[(*This Split trick pretends that two points on the
        curve are "equal" if the function values have _opposite _ sign.Pairs
        of such sign-changes form the brackets for the subsequent FindRoot*)
        Split[p, Sign[Last[#2]] == -Sign[Last[#1]] &, Length[#1] == 2 &], {2}];
    x /. Apply[FindRoot[localFunction == 0, {x, ##1}] &, brackets, {1}] /. x → {}]
```

---

Coefficients of Weyl's expansion in various settings

In[121]:=

```

αλ[λ_, μ_] := μ / (λ + 2 μ);
λα[α_, μ_] := μ (1 / α - 2);
CWeyl[d_, λ_, μ_] :=
  ((λ + 2 μ) ^ (-d / 2) + (d - 1) μ ^ (-d / 2)) / (4 Pi) ^ (d / 2) / Gamma[1 + d / 2];
BDirLiu[d_, λ_, μ_] :=
  - ((d - 1) / μ ^ ((d - 1) / 2) + 1 / (λ + 2 μ) ^ ((d - 1) / 2)) / 4 / (4 Pi) ^ ((d - 1) / 2) /
  Gamma[1 + (d - 1) / 2];
BDir[d_, α_, μ_] := -μ ^ ((1 - d) / 2) / (2 ^ (d + 1) Pi ^ ((d - 1) / 2) Gamma[(1 + d) / 2])
  (4 (d - 1) / Pi NIntegrate[τ ^ (d - 2) ArcTan[Sqrt[(1 - α τ ^ (-2)) (τ ^ (-2) - 1)]]],
  {τ, Sqrt[α], 1}] + α ^ ((d - 1) / 2) + d - 1);
BDirSaVa2[α_, μ_] := 1 / (4 Pi Sqrt[μ]) (-1 - Sqrt[α] -
  4 / Pi NIntegrate[ArcTan[Sqrt[(1 - α / ξ ^ 2) (1 / ξ ^ 2 - 1)]]], {ξ, Sqrt[α], 1}]);
BDirSaVa3[λ_, μ_] := -1 / (16 Pi) (3 λ ^ 2 + 13 λ μ + 16 μ ^ 2) / (λ ^ 2 μ + 5 λ μ ^ 2 + 6 μ ^ 3);
γR[α_] := Sqrt[Min[x /. Solve[x^3 - 8 x^2 + 16 (-1 + α) + 8 (3 - 2 α) x == 0, x, Reals]]];
Bfree[d_, α_, μ_] := μ ^ ((1 - d) / 2) / (2 ^ (d + 1) Pi ^ ((d - 1) / 2) Gamma[(1 + d) / 2])
  (4 (d - 1) / Pi NIntegrate[
  τ ^ (d - 2) ArcTan[(τ ^ (-2) - 2) ^ 2 / (4 Sqrt[(1 - α τ ^ (-2)) (τ ^ (-2) - 1)]]],
  {τ, Sqrt[α], 1}] + α ^ ((d - 1) / 2) + d - 5 + 4 γR[α] ^ (1 - d)];
BfreeSaVa2[α_, μ_] :=
  1 / (4 Pi Sqrt[μ]) (4 / γR[α] - 3 + Sqrt[α] + 4 / Pi NIntegrate[ArcTan[
  (2 - 1 / ξ ^ 2) ^ 2 / (4 Sqrt[(1 - α / ξ ^ 2) (1 / ξ ^ 2 - 1)]]], {ξ, Sqrt[α], 1}]);
BfreeSaVa3[λ_, μ_] :=
  1 / (16 Pi) (3 (λ + 2 μ) ^ 2 - 3 (λ + 2 μ) μ + 2 μ ^ 2) / ((λ + 2 μ) μ (λ + μ));

```

Lame operator,  $L$ , is the operator of plane elasticity written in polar coordinates. For vectors, the first component is along the radius.

In[132]:=

```

Lame =
  Function[u, (μ Curl[Curl[Flatten[{u, 0}], {r, φ, z}, "Cylindrical"], {r, φ, z},
  "Cylindrical"] - (λ + 2 μ) Grad[Div[Flatten[{u, 0}], {r, φ, z},
  "Cylindrical"], {r, φ, z}, "Cylindrical"])[[1 ;; 2]]];
Lame[{u1[r, φ], u2[r, φ]}]

```

Out[133]=

$$\left\{ \frac{\mu \left( -\frac{u_2^{(0,1)}[r, \phi] + u_1^{(0,2)}[r, \phi]}{r} + u_2^{(1,1)}[r, \phi] \right)}{r} - \right.$$

$$(\lambda + 2\mu) \left( -\frac{u_1[r, \phi] + u_2^{(0,1)}[r, \phi]}{r^2} + \frac{u_1^{(1,0)}[r, \phi] + u_2^{(1,1)}[r, \phi]}{r} + u_1^{(2,0)}[r, \phi] \right),$$

$$- \frac{(\lambda + 2\mu) \left( \frac{u_1^{(0,1)}[r, \phi] + u_2^{(0,2)}[r, \phi]}{r} + u_1^{(1,1)}[r, \phi] \right)}{r} +$$

$$\left. \mu \left( -\frac{-u_2[r, \phi] + u_1^{(0,1)}[r, \phi]}{r^2} + \frac{-u_2^{(1,0)}[r, \phi] + u_1^{(1,1)}[r, \phi]}{r} - u_2^{(2,0)}[r, \phi] \right) \right\}$$

The boundary conditions for the free boundary and for the Dirichlet problems are set up with

```
In[134]:=
T[u_] := ({λ Div[u, {r, φ}, "Polar"], 0} + 2 μ D[u, r] +
  μ ({1, 0, 0} × Curl[Flatten[{u, 0}], {r, φ, z}, "Cylindrical"]) [[1 ;; 2]] /.
  {r → 1}) // Simplify;
TDir[u_] := {u[[1]], u[[2]]};
```

```
In[136]:=
T[{u1[r, φ], u2[r, φ]}]
TDir[{u1[r, φ], u2[r, φ]}]
```

```
Out[136]=
{λ u1[1, φ] + λ u2(0,1)[1, φ] + (λ + 2 μ) u1(1,0)[1, φ],
  μ (-u2[1, φ] + u1(0,1)[1, φ] + u2(1,0)[1, φ])}
```

```
Out[137]=
{u1[r, φ], u2[r, φ]}
```

---

Solutions of Lamé system for  $\Lambda=0$ : there are 4 solutions,  $\{r,0\}$ ,  $\{0,r\}$ ,  $\{1, l\} \text{Exp}[l \phi]$ ,  $\{1, -l\} \text{Exp}[-l \phi]$   
 This is irrelevant for the discussion and can be skipped

```
In[138]:=
Collect[Lame[{a r, c r}], r, Simplify]
T[{a r, c r}]
```

```
Out[138]=
{0, 0}
```

```
Out[139]=
{2 a (λ + μ), 0}
```

```
In[140]:=
Collect[Exp[-I φ] Lamé[{b, I b} Exp[I φ]], r, Simplify]
T[{b, I b} Exp[I φ]]
```

```
Out[140]=
{0, 0}
```

```
Out[141]=
{0, 0}
```

```
In[142]:=
Collect[Exp[I φ] Lamé[{b, -I b} Exp[-I φ]], r, Simplify]
T[{b, -I b} Exp[-I φ]]
```

```
Out[142]=
{0, 0}
```

```
Out[143]=
{0, 0}
```

```
In[144]:=
κ0 = (λ + 3 μ) / (λ + μ)
```

```
Out[144]=

$$\frac{\lambda + 3 \mu}{\lambda + \mu}$$

```

In[145]:=

```
ClearAll[soln0]; soln0[n_] := Which[
  n == 0, {{r, 0}, {0, r}},
  n == 1, {{1, I}, r^2 {(x0 - 2), -I (x0 + 2)}},
  n == -1, {{1, -I}, r^2 {(x0 - 2), I (x0 + 2)}},
  n ≥ 2, {n r^ (n - 1) {1, I}, r^ (n + 1) {(x0 - n - 1), -I (x0 + n + 1)}},
  n ≤ -2, {Abs[n] r^ (Abs[n] - 1) {1, -I},
  r^ (Abs[n] + 1) {(x0 - Abs[n] - 1), I (x0 + Abs[n] + 1)}}];
```

Formula for  $|n| \geq 2$  in fact works for  $|n|=1$ , and we therefore redefine soln0:

In[146]:=

```
{(n r^ (n - 1) {1, I}, r^ (n + 1) {(x0 - n - 1), -I (x0 + n + 1)}) /. n -> 1} -
  {{1, I}, r^2 {(x0 - 2), -I (x0 + 2)}}
```

Out[146]=

```
{{0, 0}, {0, 0}}
```

In[148]:=

```
soln0[n_] := Piecewise[{{
  {{r, 0}, {0, r}}, n == 0,
  {{n r^ (n - 1) {1, I}, r^ (n + 1) {(x0 - n - 1), -I (x0 + n + 1)}} Exp[I n φ], n ≥ 1,
  {{Abs[n] r^ (Abs[n] - 1) {1, -I},
  r^ (Abs[n] + 1) {(x0 - Abs[n] - 1), I (x0 + Abs[n] + 1)}} Exp[I n φ], n ≤ -1
}}];
```

In[149]:=

```
Table[Map[Lame, soln0[n]], {n, -4, 4}] // Simplify
```

Out[149]=

```
{{{0, 0}, {0, 0}}, {{0, 0}, {0, 0}}, {{0, 0}, {0, 0}},
  {{0, 0}, {0, 0}}, {{0, 0}, {0, 0}}, {{0, 0}, {0, 0}},
  {{0, 0}, {0, 0}}, {{0, 0}, {0, 0}}, {{0, 0}, {0, 0}}
```

In[150]:=

```
Table[Map[T, soln0[n]], {n, -2, 2}] // Simplify
```

Out[150]=

```
{{{4 e^{-2 i φ} μ, -4 i e^{-2 i φ} μ}, {0, 12 i e^{-2 i φ} μ}},
  {{0, 0}, {4 e^{-i φ} μ, 4 i e^{-i φ} μ}}, {{2 (λ + μ), 0}, {0, 0}},
  {{0, 0}, {4 e^{i φ} μ, -4 i e^{i φ} μ}}, {{4 e^{2 i φ} μ, 4 i e^{2 i φ} μ}, {0, -12 i e^{2 i φ} μ}}}
```

In[151]:=

```
Table[Map[TDir, soln0[n]], {n, -2, 2}] // Simplify
```

Out[151]=

```
{{{2 e^{-2 i φ} r, -2 i e^{-2 i φ} r}, {-\frac{2 e^{-2 i φ} r^3 λ}{λ + μ}, \frac{2 i e^{-2 i φ} r^3 (2 λ + 3 μ)}{λ + μ}}},
  {{e^{-i φ}, -i e^{-i φ}}, {-\frac{e^{-i φ} r^2 (λ - μ)}{λ + μ}, \frac{i e^{-i φ} r^2 (3 λ + 5 μ)}{λ + μ}}},
  {{r, 0}, {0, r}}, {{e^{i φ}, i e^{i φ}}, {-\frac{e^{i φ} r^2 (λ - μ)}{λ + μ}, -\frac{i e^{i φ} r^2 (3 λ + 5 μ)}{λ + μ}}},
  {{2 e^{2 i φ} r, 2 i e^{2 i φ} r}, {-\frac{2 e^{2 i φ} r^3 λ}{λ + μ}, -\frac{2 i e^{2 i φ} r^3 (2 λ + 3 μ)}{λ + μ}}}
```

Solutions of the Lamé system for  $\Lambda \neq 0$ .

We seek these solutions in the form  $\text{Grad}[\Phi[r, \phi]] + \text{Curl}[\vec{Z}\Psi[r, \phi]]$ . Then the scalar potentials  $\Phi$

and  $\Psi$  should satisfy Helmholtz equations, as seen from the calculation below. Indeed, setting  $u_g = \text{Grad}[\Phi[r, \phi]]$  and  $u_c = \text{Curl}[\vec{z}\Psi[r, \phi]]$ , we get

$$\text{Lame}[u_g] - \Delta u_g = \text{Grad}[-(\lambda+2\mu)\Delta\Phi - \Lambda\Phi]$$

and

$$\text{Lame}[u_c] - \Delta u_c = \text{Curl}[\vec{z}(-\mu\Delta\Psi - \Lambda\Psi)]$$

In[152]:=

```
Curl[Curl[{0, 0, Ψ[r, ϕ]}, {r, ϕ, z}, "Cylindrical"],
      {r, ϕ, z}, "Cylindrical"][[3]]
Laplacian[Ψ[r, ϕ], {r, ϕ, z}, "Cylindrical"]
```

Out[152]=

$$-\frac{\frac{\Psi^{(0,2)}[r, \phi]}{r} + \Psi^{(1,0)}[r, \phi]}{r} - \Psi^{(2,0)}[r, \phi]$$

Out[153]=

$$\frac{\frac{\Psi^{(0,2)}[r, \phi]}{r} + \Psi^{(1,0)}[r, \phi]}{r} + \Psi^{(2,0)}[r, \phi]$$

In[154]:=

```
ug = Grad[Ψ[r, ϕ], {r, ϕ, z}, "Cylindrical"][[1 ;; 2]]
eqg1 = FullSimplify[Lame[ug] - Δug];
eqg2 = FullSimplify[
  Grad[-(λ+2μ) Laplacian[Ψ[r, ϕ], {r, ϕ, z}, "Cylindrical"] -
    ΔΨ[r, ϕ], {r, ϕ, z}, "Cylindrical"][[1 ;; 2]];
eqg1 - eqg2
```

Out[154]=

$$\left\{ \Psi^{(1,0)}[r, \phi], \frac{\Psi^{(0,1)}[r, \phi]}{r} \right\}$$

Out[157]=

$$\{0, 0\}$$

In[158]:=

```
uc = Curl[{0, 0, Ψ[r, ϕ]}, {r, ϕ, z}, "Cylindrical"][[1 ;; 2]]
eqc1 = FullSimplify[Lame[uc] - Δuc];
eqc2 = FullSimplify[
  Curl[{0, 0, -μ Laplacian[Ψ[r, ϕ], {r, ϕ, z}, "Cylindrical"] - ΔΨ[r, ϕ]},
    {r, ϕ, z}, "Cylindrical"][[1 ;; 2]];
eqc1 - eqc2 // Simplify
```

Out[158]=

$$\left\{ \frac{\Psi^{(0,1)}[r, \phi]}{r}, -\Psi^{(1,0)}[r, \phi] \right\}$$

Out[161]=

$$\{0, 0\}$$

The solutions of Helmholtz equation  $-\Delta\theta - k\theta = 0$  are

$$\theta = J_{|n|}(\sqrt{k} r) \exp(in\phi) \text{ for } k > 0 \text{ and } \theta = I_{|n|}(\sqrt{-k} r) \exp(in\phi) \text{ for } k < 0.$$

We will only deal with non-negative  $n$ 's remembering later to double the multiplicities for  $n \neq 0$ . `solnneg` refers to  $k < 0$  and may be ignored

In[162]:=

```
solnpos[n_] :=
  {Grad[BesselJ[n, Sqrt[ $\Lambda / (\lambda + 2 \mu)$ ] r] Exp[I n  $\phi$ ], {r,  $\phi$ , z}, "Cylindrical"][[1 ;; 2]],
  Curl[{0, 0, BesselJ[n, Sqrt[ $\Lambda / \mu$ ] r] Exp[I n  $\phi$ ]], {r,  $\phi$ , z}, "Cylindrical"][[
  1 ;; 2]] /. BesselJ[-1 + n, z_]  $\rightarrow$  (2 n BesselJ[n, z] / z - BesselJ[n + 1, z])
solnneg[n_] :=
  {Grad[BesselI[n, Sqrt[- $\Lambda / (\lambda + 2 \mu)$ ] r] Exp[I n  $\phi$ ], {r,  $\phi$ , z}, "Cylindrical"][[
  1 ;; 2]], Curl[{0, 0, BesselI[n, Sqrt[- $\Lambda / \mu$ ] r] Exp[I n  $\phi$ ]],
  {r,  $\phi$ , z}, "Cylindrical"][[1 ;; 2]] /.
  BesselI[-1 + n, z_]  $\rightarrow$  (2 n BesselI[n, z] / z + BesselI[n + 1, z]) // FullSimplify;
```

Just to check that we indeed produce solutions of Lamé system for both positive and negative n:

In[164]:=

```
Map[(Lame[#] -  $\Lambda$  #) &, solnpos[n]] // FullSimplify
```

Out[164]=

```
{{0, 0}, {0, 0}}
```

In[165]:=

```
solnposMinusn[n_] :=
  {Grad[BesselJ[n, Sqrt[ $\Lambda / (\lambda + 2 \mu)$ ] r] Exp[-I n  $\phi$ ], {r,  $\phi$ , z}, "Cylindrical"][[
  1 ;; 2]], Curl[{0, 0, BesselJ[n, Sqrt[ $\Lambda / \mu$ ] r] Exp[-I n  $\phi$ ]],
  {r,  $\phi$ , z}, "Cylindrical"][[1 ;; 2]] /.
  BesselJ[-1 + n, z_]  $\rightarrow$  (2 n BesselJ[n, z] / z - BesselJ[n + 1, z]);
```

In[166]:=

```
Map[(Lame[#] -  $\Lambda$  #) &, solnposMinusn[n]] // FullSimplify
```

Out[166]=

```
{{0, 0}, {0, 0}}
```

This is to check that choosing plus or minus n does not affect b.c. : the first (gradient) solutions changes the radial part as  $(u,v) \rightarrow (u,-v)$  under  $n \rightarrow -n$ , the second (curl) solution changes the radial part as  $(u,v) \rightarrow (-u,v)$  under  $n \rightarrow -n$ , so homogeneous conditions stay the same. So from now on we just deal with  $n \geq 0$  and double all the eigenvalues for non-zero n

In[167]:=

```

Exp[-I n ϕ] solnpos[n][[1]] -
  Exp[I n ϕ] {{1, 0}, {0, -1}} . solnposMinusn[n][[1]] // Simplify
Exp[-I n ϕ] solnpos[n][[2]] -
  Exp[I n ϕ] {{-1, 0}, {0, 1}} . solnposMinusn[n][[2]] // Simplify
Exp[-I n ϕ] T[solnpos[n][[1]]] -
  Exp[I n ϕ] {{1, 0}, {0, -1}} . T[solnposMinusn[n][[1]]] // Simplify
Exp[-I n ϕ] T[solnpos[n][[2]]] -
  Exp[I n ϕ] {{-1, 0}, {0, 1}} . T[solnposMinusn[n][[2]]] // Simplify

```

Out[167]=

```
{0, 0}
```

Out[168]=

```
{0, 0}
```

Out[169]=

```
{0, 0}
```

Out[170]=

```
{0, 0}
```

The Dirichlet eigenvalues are the  $\Lambda$ -roots of

In[171]:=

```
(Map[#, {r → 1, Exp[_] → 1}] &, solnpos[n]] // Det // FullSimplify
```

Out[171]=

$$\begin{aligned}
& n \sqrt{\frac{\Lambda}{\mu}} \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\lambda + 2\mu}}\right] \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\mu}}\right] + \\
& \sqrt{\frac{\Lambda}{\lambda + 2\mu}} \left( n \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\mu}}\right] - \sqrt{\frac{\Lambda}{\mu}} \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\mu}}\right] \right) \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\lambda + 2\mu}}\right]
\end{aligned}$$

In[172]:=

```

Direigseq = (Map[#, {r → 1, Exp[_] → 1}] &, solnpos[n]] // Det /.
  BesselJ[-1 + n, z_] → (2 n BesselJ[n, z] / z - BesselJ[n + 1, z]) // FullSimplify

```

Out[172]=

$$\begin{aligned}
& n \sqrt{\frac{\Lambda}{\mu}} \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\lambda + 2\mu}}\right] \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\mu}}\right] + \\
& \sqrt{\frac{\Lambda}{\lambda + 2\mu}} \left( n \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\mu}}\right] - \sqrt{\frac{\Lambda}{\mu}} \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\mu}}\right] \right) \text{BesselJ}\left[1 + n, \sqrt{\frac{\Lambda}{\lambda + 2\mu}}\right]
\end{aligned}$$

In[183]:=

$$\text{Direigseq}\omega = \left( \text{Direigseq} /. \left\{ \frac{\Lambda}{\lambda + 2\mu} \rightarrow \omega_{1,\Lambda}, \frac{\Lambda}{\mu} \rightarrow \omega_{2,\Lambda}, n \rightarrow k \right\} \right)$$

Direigseq  $\omega$  /.  $n \rightarrow 0$ Direigseq  $\omega$  // TeXForm

Out[183]=

$$\text{BesselJ}\left[1+k, \sqrt{\omega_{1,\Lambda}}\right] \sqrt{\omega_{1,\Lambda}} \left( k \text{BesselJ}\left[k, \sqrt{\omega_{2,\Lambda}}\right] - \text{BesselJ}\left[1+k, \sqrt{\omega_{2,\Lambda}}\right] \sqrt{\omega_{2,\Lambda}} \right) + k \text{BesselJ}\left[k, \sqrt{\omega_{1,\Lambda}}\right] \text{BesselJ}\left[1+k, \sqrt{\omega_{2,\Lambda}}\right] \sqrt{\omega_{2,\Lambda}}$$

Out[184]=

$$\text{BesselJ}\left[1+k, \sqrt{\omega_{1,\Lambda}}\right] \sqrt{\omega_{1,\Lambda}} \left( k \text{BesselJ}\left[k, \sqrt{\omega_{2,\Lambda}}\right] - \text{BesselJ}\left[1+k, \sqrt{\omega_{2,\Lambda}}\right] \sqrt{\omega_{2,\Lambda}} \right) + k \text{BesselJ}\left[k, \sqrt{\omega_{1,\Lambda}}\right] \text{BesselJ}\left[1+k, \sqrt{\omega_{2,\Lambda}}\right] \sqrt{\omega_{2,\Lambda}}$$

Out[185]//TeXForm=

$$k \sqrt{\omega_{2,\Lambda}} J_k\left(\sqrt{\omega_{1,\Lambda}}\right) J_{k+1}\left(\sqrt{\omega_{2,\Lambda}}\right) + \sqrt{\omega_{1,\Lambda}} J_{k+1}\left(\sqrt{\omega_{1,\Lambda}}\right) \left( k J_k\left(\sqrt{\omega_{2,\Lambda}}\right) - \sqrt{\omega_{2,\Lambda}} J_{k+1}\left(\sqrt{\omega_{2,\Lambda}}\right) \right) - \sqrt{\omega_{2,\Lambda}} J_{k+1}\left(\sqrt{\omega_{1,\Lambda}}\right) \left( k J_k\left(\sqrt{\omega_{2,\Lambda}}\right) - \sqrt{\omega_{2,\Lambda}} J_{k+1}\left(\sqrt{\omega_{2,\Lambda}}\right) \right)$$

and the Neumann eigenvalues are  $\Lambda$ -roots of (this sometimes does not show as Bessels but as Hypergeometric functions)

In[186]:=

Neueigseq = (Map[(T[#] /. {r → 1, Exp[\_] → 1}) &amp;, solnpos[n]] // Det) /.

BesselJ[-1+n, z] → (2 n BesselJ[n, z] / z - BesselJ[n+1, z]) // FullSimplify

Out[186]=

$$2 \sqrt{\frac{\Lambda}{\mu}} \mu \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\mu}}\right] \left( (\Lambda + 2 n \mu - 2 n^3 \mu) \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right] + 2 (-1 + n^2) \mu \sqrt{\frac{\Lambda}{\lambda + 2 \mu}} \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right] \right) + \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\mu}}\right] \left( -\Lambda (\Lambda - 4 (-1 + n) n \mu) \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right] + 2 \mu \sqrt{\frac{\Lambda}{\lambda + 2 \mu}} (\Lambda + 2 n \mu - 2 n^3 \mu) \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right] \right)$$

In[192]:=

Direq[ $\lambda_$ ,  $\mu_$ ,  $n_$ ,  $\Lambda_$ ] :=

$$n \sqrt{\frac{1}{\mu}} \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right] \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\mu}}\right] + \sqrt{\frac{1}{\lambda + 2 \mu}} \left( n \text{BesselJ}\left[n, \sqrt{\frac{\Lambda}{\mu}}\right] - \sqrt{\frac{\Lambda}{\mu}} \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\mu}}\right] \right) \text{BesselJ}\left[1+n, \sqrt{\frac{\Lambda}{\lambda + 2 \mu}}\right];$$

Simplify[Direq[ $\lambda$ ,  $\mu$ ,  $n$ ,  $\Lambda$ ] - Direigseq / Sqrt[ $\Lambda$ ],  $\Lambda > 0 \ \&\& \ \mu > 0 \ \&\& \ \lambda + 2 \mu > 0$ ]

Out[196]=

0



In[197]:=

**Direq**[ $\lambda, \mu, n, \Delta$ ] // FullSimplify  
**Direq**[ $\lambda, \mu, 0, \Delta$ ] // Simplify

Out[197]=

$$n \sqrt{\frac{1}{\mu}} \text{BesselJ}\left[n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right] \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\mu}}\right] +$$

$$\sqrt{\frac{1}{\lambda + 2\mu}} \left( n \text{BesselJ}\left[n, \sqrt{\frac{\Delta}{\mu}}\right] - \sqrt{\frac{\Delta}{\mu}} \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\mu}}\right] \right) \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right]$$

Out[198]=

$$- \sqrt{\frac{\Delta}{\mu}} \sqrt{\frac{1}{\lambda + 2\mu}} \text{BesselJ}\left[1, \sqrt{\frac{\Delta}{\mu}}\right] \text{BesselJ}\left[1, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right]$$

**Neueq**[ $\lambda_-, \mu_-, n_-, \Delta_-$ ] :=

$$2 \sqrt{\frac{\Delta}{\mu}} \mu \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\mu}}\right] \left( (\Delta + 2 n \mu - 2 n^3 \mu) \text{BesselJ}\left[n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right] + \right.$$

$$2 (-1 + n^2) \mu \sqrt{\frac{\Delta}{\lambda + 2\mu}} \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right] \left. \right) +$$

$$\text{BesselJ}\left[n, \sqrt{\frac{\Delta}{\mu}}\right] \left( -\Delta (\Delta - 4 (-1 + n) n \mu) \text{BesselJ}\left[n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right] + \right.$$

$$2 \mu \sqrt{\frac{\Delta}{\lambda + 2\mu}} (\Delta + 2 n \mu - 2 n^3 \mu) \text{BesselJ}\left[1 + n, \sqrt{\frac{\Delta}{\lambda + 2\mu}}\right] \left. \right);$$

**Simplify**[**Neueq**[ $\lambda, \mu, n, \Delta$ ] - **Neueigseq**,  $\Delta > 0 \&\& \mu > 0 \&\& \lambda + 2 \mu > 0$ ]

Out[200]=

0

In[201]:=

```

Direigs[λ0_, μ0_, nmax_, Δmax_] := Module[{t, j},
  t = Table[findAllRoots[Direq[λ0, μ0, n, Δ], {Δ, 0., Δmax}], {n, 0, nmax}];
  t = {t[[1]], Table[{t[[j]], t[[j]]}, {j, 2, Length[t]}]} // Flatten;
  t // Sort
];
DireigsWithn[λ0_, μ0_, nmax_, Δmax_] := Module[{t, t1, t2, t3, pair},
  t = Table[findAllRoots[Direq[λ0, μ0, n, Δ], {Δ, 0., Δmax}], {n, 0, nmax}];
  t1 = Select[Table[If[Length[t[[n + 1]]] > 0, Thread[{n, t[[n + 1]]}],
    {n, 0, nmax}], Length[#] > 0 &];
  t2 = SortBy[Flatten[t1, 1], N[#[[2]]] &];
  t2
];
Neueigs[λ0_, μ0_, nmax_, Δmax_] := Module[{t, k},
  t = Table[findAllRoots[Neueq[λ0, μ0, n, Δ], {Δ, 0., Δmax}], {n, 0, nmax}];
  t = {t[[1]], Table[{t[[j]], t[[j]]}, {j, 2, Length[t]}]} // Flatten;
  {0., 0., 0., t // Sort} // Flatten
];
NeueigsWithn[λ0_, μ0_, nmax_, Δmax_] := Module[{t, t1, t2, t3, pair},
  t = Table[findAllRoots[Neueq[λ0, μ0, n, Δ], {Δ, 0., Δmax}], {n, 0, nmax}];
  t1 = Select[Table[If[Length[t[[n + 1]]] > 0, Thread[{n, t[[n + 1]]}],
    {n, 0, nmax}], Length[#] > 0 &];
  t2 = SortBy[Flatten[t1, 1], #[[2]] &];
  PrependTo[t2, {1, 0.}];
  PrependTo[t2, {0, 0.}];
  t2
];

```

Examples:

In[205]:=

```

Direigs[0, 1, 52, 3200.]
DireigsWithn[0, 1, 52, 3200.]

```

Out[205]=

```

{8.6126, 8.6126, 14.682, 21.5228, 21.5228, 28.4568, 28.4568, 29.3639, 37.877,
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Out[206]=

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{18, 2997.87}, {17, 3010.35}, {36, 3010.45}, {24, 3015.23}, {21, 3016.95},
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{6, 3061.55}, {46, 3066.28}, {41, 3066.72}, {26, 3066.99}, {20, 3068.57},
{31, 3070.53}, {4, 3086.95}, {15, 3091.6}, {23, 3093.56}, {8, 3094.6},
{2, 3101.18}, {0, 3105.79}, {30, 3120.05}, {16, 3125.23}, {11, 3134.23},
{6, 3140.14}, {28, 3144.04}, {34, 3148.94}, {37, 3155.53}, {13, 3165.82},
{19, 3166.77}, {25, 3170.01}, {18, 3172.21}, {4, 3175.65}, {38, 3176.8},
{22, 3179.65}, {47, 3183.65}, {42, 3188.67}, {9, 3189.9}, {2, 3197.84}}
```

In[209]:=

```
MyNlambdaDisk[evs_,  $\Lambda$ ] :=
  Sum[If[evs[[j]][[1]] == 0, 1, 2] UnitStep[ $\Lambda$  - evs[[j]][[2]]], {j, 1, Length[evs]}];
```

Finding the number of harmonics to stop at

In[210]:=

```
findAllRoots[Direq[-0.5, 1, 47,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
findAllRoots[Direq[-0.5, 1, 48,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
```

Out[210]=

```
{2961.62}
```

Out[211]=

```
{}
```

In[212]:=

```
findAllRoots[Direq[3, 1, 46,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
findAllRoots[Direq[3, 1, 47,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
```

Out[212]=

```
{2900.41}
```

Out[213]=

```
{}
```

In[214]:=

```
findAllRoots[Direq[100, 1, 46,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
findAllRoots[Direq[100, 1, 47,  $\Lambda$ ], { $\Lambda$ , 0., 3000}]
```

Out[214]=

```
{2913.57}
```

Out[215]=

```
{}
```

Dirichlet counting functions

In[216]:=

```
evD1 = DireigsWithn[-0.5, 1, 47, 3000];
evD2 = DireigsWithn[3, 1, 46, 3000];
evD3 = DireigsWithn[100, 1, 46, 3000];
```

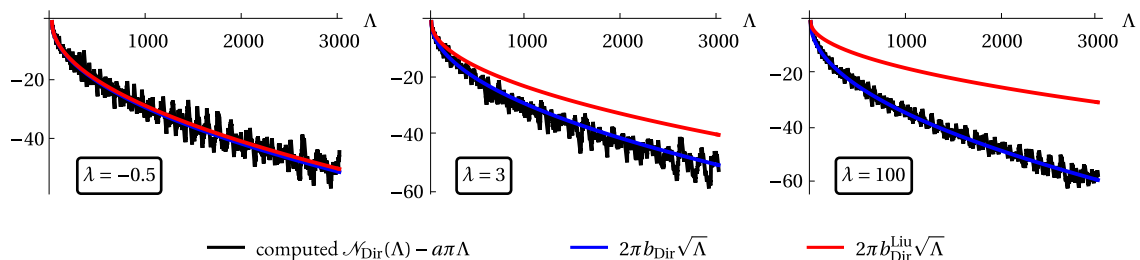
In[219]:=

```
xtdisk = Table[{x, MaTeX[x]}, {x, 1000, 3000, 1000}];
ytdisk = Table[{x, MaTeX[x]}, {x, -60, -20, 20}];
```

In[221]:=

```
figdisk1 = Plot[{MyNlambdaDisk[evD1, Δ] - CWeyl[2, -0.5, 1] Pi Δ,
  2 Pi BDir[2, αλ[-0.5, 1], 1] Sqrt[Δ], 2 Pi BDirLiu[2, -0.5, 1] Sqrt[Δ]},
  {Δ, 0, 3000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytdisk}, AxesLabel → {MaTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MaTeX["\\lambda=-0.5"], RoundingRadius → 2, ContentPadding →
  False, FrameMargins → Tiny], Scaled[{0.1, 0}], Scaled[{0, 0}]]];
figdisk2 = Plot[{MyNlambdaDisk[evD2, Δ] - CWeyl[2, 3, 1] Pi Δ,
  2 Pi BDir[2, αλ[3, 1], 1] Sqrt[Δ], 2 Pi BDirLiu[2, 3, 1] Sqrt[Δ]},
  {Δ, 0, 3000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytdisk}, AxesLabel → {MaTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MaTeX["\\lambda=3"], RoundingRadius → 2, ContentPadding →
  False, FrameMargins → Tiny], Scaled[{0.1, 0}], Scaled[{0, 0}]]];
figdisk3 = Plot[{MyNlambdaDisk[evD3, Δ] - CWeyl[2, 100, 1] Pi Δ,
  2 Pi BDir[2, αλ[100, 1], 1] Sqrt[Δ], 2 Pi BDirLiu[2, 100, 1] Sqrt[Δ]},
  {Δ, 0, 3000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytdisk}, AxesLabel → {MaTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MaTeX["\\lambda=100"], RoundingRadius → 2, ContentPadding →
  False, FrameMargins → Tiny], Scaled[{0.1, 0}], Scaled[{0, 0}]]];
legendisk = LineLegend[{Black, Blue, Red},
  MaTeX[{"\\text{computed } \\mathcal{N}_{\\text{Dir}}(\\Lambda) - a\\pi \\Lambda",
  "\\Lambda\\quad\\quad\\quad",
  "2\\pi b_{\\text{Dir}}\\sqrt{\\Lambda}\\quad\\quad\\quad",
  "2\\pi b_{\\text{Dir}}^{\\text{Liu}}\\sqrt{\\Lambda}}"],
  LegendFunction → None, LegendLayout → {"Row", 1}];
figdisk = GraphicsColumn[{GraphicsRow[{figdisk1, figdisk2, figdisk3},
  ImageSize → Full, Spacings → {Scaled[0.05], Automatic}],
  legendisk}, ImageSize → Full, Spacings → {0, Automatic}]
```

Out[225]=



In[\*]:= Export["figdisk.pdf", figdisk]

Out[\*]=

figdisk.pdf

Finding the number of harmonics to stop at

```
In[226]:=
  findAllRoots[Neueq[-0.5, 1, 56,  $\Delta$ ], { $\Delta$ , 0., 2000}]
  findAllRoots[Neueq[-0.5, 1, 57,  $\Delta$ ], { $\Delta$ , 0., 2000}]
```

```
Out[226]=
  {1992.}
```

```
Out[227]=
  {}
```

```
In[228]:=
  findAllRoots[Neueq[3, 1, 45,  $\Delta$ ], { $\Delta$ , 0., 2000}]
  findAllRoots[Neueq[3, 1, 46,  $\Delta$ ], { $\Delta$ , 0., 2000}]
```

```
Out[228]=
  {1929.28}
```

```
Out[229]=
  {}
```

```
In[230]:=
  findAllRoots[Neueq[100, 1, 44,  $\Delta$ ], { $\Delta$ , 0., 2000}]
  findAllRoots[Neueq[100, 1, 45,  $\Delta$ ], { $\Delta$ , 0., 2000}]
```

```
Out[230]=
  {1915.88}
```

```
Out[231]=
  {}
```

---

```
In[232]:=
  evDfr1 = NeueigsWithn[-0.5, 1, 56, 2000];
  evDfr2 = NeueigsWithn[3, 1, 45, 2000];
  evDfr3 = NeueigsWithn[100, 1, 44, 2000];
```

```
In[235]:=
  ytfdisk = Table[{x, MaTeX["\\ " <> ToString[x]]}, {x, 20, 60, 20}];
```





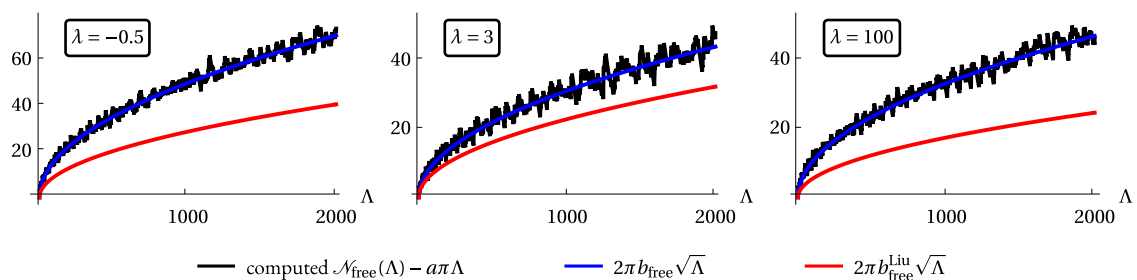
In[241]=

```

figdiskfr1new = Plot[{MyNlambdaDisk[evDfr1, Δ] - CWeyl[2, -0.5, 1] Pi Δ,
  2 Pi Bfree[2, αλ[-0.5, 1], 1] Sqrt[Δ], -2 Pi BDirLiu[2, -0.5, 1] Sqrt[Δ]},
  {Δ, 0, 2000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytfrdisk}, AxesLabel → {MathTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MathTeX["\\lambda=-0.5"], RoundingRadius → 2, ContentPadding →
    False, FrameMargins → Tiny], Scaled[{0.1, 1}], Scaled[{0, 1}]]];
figdiskfr2new = Plot[{MyNlambdaDisk[evDfr2, Δ] - CWeyl[2, 3, 1] Pi Δ,
  2 Pi Bfree[2, αλ[3, 1], 1] Sqrt[Δ], -2 Pi BDirLiu[2, 3, 1] Sqrt[Δ]},
  {Δ, 0, 2000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytfrdisk}, AxesLabel → {MathTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MathTeX["\\lambda=3"], RoundingRadius → 2, ContentPadding →
    False, FrameMargins → Tiny], Scaled[{0.1, 1}], Scaled[{0, 1}]]];
figdiskfr3new = Plot[{MyNlambdaDisk[evDfr3, Δ] - CWeyl[2, 100, 1] Pi Δ,
  2 Pi Bfree[2, αλ[100, 1], 1] Sqrt[Δ], -2 Pi BDirLiu[2, 100, 1] Sqrt[Δ]},
  {Δ, 0, 2000}, PlotStyle → {Black, Blue, Red}, Exclusions → None,
  Ticks → {xtdisk, ytfrdisk}, AxesLabel → {MathTeX["\\Lambda"], None}, Epilog →
  Inset[Framed[MathTeX["\\lambda=100"], RoundingRadius → 2, ContentPadding →
    False, FrameMargins → Tiny], Scaled[{0.1, 1}], Scaled[{0, 1}]]];
legdiskfrnew = LineLegend[{Black, Blue, Red},
  MathTeX[{"\\text{computed } \\mathcal{N}_{\\mathrm{free}}(\\Lambda) - a\\pi \\Lambda \\quad \\quad \\quad",
    "2\\pi b_{\\mathrm{free}} \\sqrt{\\Lambda} \\quad \\quad \\quad",
    "2\\pi b_{\\mathrm{free}}^{\\mathrm{Liu}} \\sqrt{\\Lambda}"}],
  LegendFunction → None, LegendLayout → {"Row", 1}];
figdiskfrnew =
GraphicsColumn[{GraphicsRow[{figdiskfr1new, figdiskfr2new, figdiskfr3new},
  ImageSize → Full, Spacings → {Scaled[0.05], Automatic}],
  legdiskfrnew}, ImageSize → Full, Spacings → {0, Automatic}]

```

Out[245]=



In[ ]:= Export["figdiskfreenew.pdf", figdiskfrnew]

Out[ ]:=

figdiskfreenew.pdf